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Illumination of mesospheric irregularity by lightning discharge

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[1] Theoretical model calculations recently predicted the existence of mesospheric irregularities which assist the initiation of sprites. Here we report the experimental detection of a $\sim 3\text{--}19\text{ km}^3$ large mesospheric irregularity at $\sim 80\text{--}85\text{ km}$ height which is illuminated by the electromagnetic field of an intense positive cloud-to-ground lightning discharge. While the lightning discharge causes a prompt group of four sprites above the lightning discharge, the mesospheric irregularity is found at a horizontal distance at least $\sim 15\text{--}20\text{ km}$ away from the sprite group and it rebrightens $\sim 40\text{--}60\text{ ms}$ after the sprite group occurrence. This rebrightening is driven by a local quasi-static electric field enhancement with a charge moment $\sim 4\text{--}20\text{ C km}$ which causes the irregularity to develop a downward descending luminous column from $\sim 75\text{--}85\text{ km}$ height. The quasi-static electric field enhancement is caused by the reorganization of residual charge inside the thundercloud during a high-level activity of intracloud discharges with $\sim 10\text{--}20$ pulses per ms. Such mesospheric irregularities might have an effect on the wave propagation of 100 kHz radio waves which are used for atomic time transfer and marine navigation. **Citation:** Füllekrug, M., A. Mezentsev, S. Soula, O. van der Velde, and A. Evans (2013), Illumination of mesospheric irregularity by lightning discharge, *Geophys. Res. Lett.*, 40, 6411–6416, doi:10.1002/2013GL058502.

1. Introduction

[2] Recent model calculations strongly suggest that enhancements of the electron density in the middle atmosphere are particularly effective in assisting the initiation of sprites below the conventional breakdown threshold [Qin *et al.*, 2013b, 2013a, 2012a, 2012b; Kosar *et al.*, 2012]. The initiation of sprites below the conventional breakdown threshold is currently under debate [Liu *et al.*, 2012] based on numerous compelling observations and numerical simulations of sprites [Gamerota *et al.*, 2011; Li *et al.*, 2008; Hu *et al.*, 2007]. The prediction that inhomogeneities in the

middle atmosphere could be illuminated by lightning electromagnetic fields is a major milestone for the maturing understanding of sprites because the physical mechanisms leading to sprite streamers might be used as a remote-sensing tool for the middle atmosphere in future studies. For example, sprites have already been used to infer the electric field strengths produced by a preceding gigantic jet [Neubert *et al.*, 2011]. Moreover, electron density inhomogeneities in the mesosphere might have a significant impact on the absorption and scattering of radio waves which will be used for the remote sensing of lightning discharges and sprites on the forthcoming TARANIS (Tool for the Analysis of RAdiations from lightNings and Sprites) satellite mission [Lefeuvre *et al.*, 2008; Blanc *et al.*, 2007]. An improved understanding of the proposed inhomogeneities in the middle atmosphere is therefore important for the preparation of the TARANIS mission, which is the aim of this contribution.

2. Observations

[3] Unstable air masses between France and Spain developed into a mesoscale convective system in the early evening hours of 16 August 2013 (Figure 1). The convective system reaches an extent of $\sim 300 \times 300\text{ km}^2$ around $\sim 20\text{--}21\text{ UTC}$ and covers the entire mountain range of the Pyrénées where numerous sprite occurrences have been predicted to occur [Füllekrug and Reising, 1998]. A positive lightning discharge with a peak current of $\sim 159\text{ kA}$ occurs at 20:27:58.721 UTC. Several other weak negative lightning discharges occur at scattered locations up to $\pm 400\text{ ms}$ before and after the positive lightning discharge (not shown in Figure 1). The area above the mesoscale convective system is observed with a low-light video camera located in Clermont-Ferrand (45.8°N , 3.1°E , 400 m), $\sim 440\text{ km}$ north-east of the positive lightning discharge. This low-light video camera is a Wattec 902H, which records video frames with 8 bit resolution using a fixed nonlinear (gamma) gain and stepped interlacing at an initial sampling frequency of 25 frames per second [Soula *et al.*, 2012]. The subsequent deinterlacing into even and odd fields during post-processing results in an image sequence with a temporal resolution of 20 ms as verified by meteor track observations. The camera records a prompt group of four carrot sprites (Figure 2a) in the image lasting from 20:27:58.718–58.738, which is synchronized to UTC with a GPS time inserter with an accuracy $\sim 1\text{ ms}$. The sprite group occurs around the azimuth of the positive lightning discharge (Figure 1) as inferred from the video image with an accuracy $\sim 1\text{ km}$ [van der Velde *et al.*, 2010] by use of the software “Cartes du Ciel” (sky charts), which matches stars with those visible in the video image. The horizontal extent of the sprite group is estimated to be $\sim 30\text{--}40\text{ km}$ by using the distance between the camera and the positive lightning discharge.

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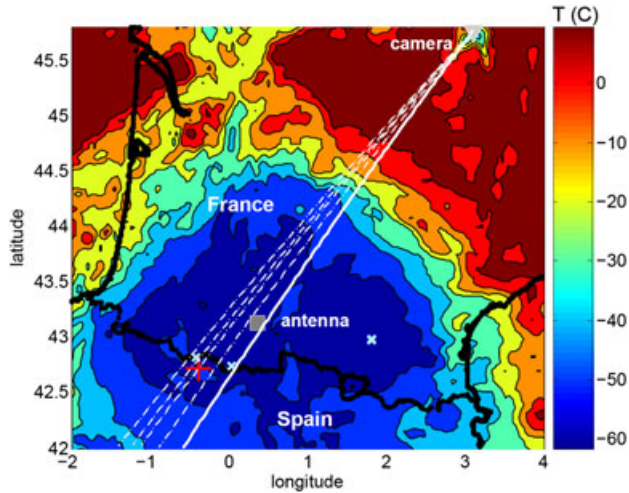


Figure 1. A mesoscale convective system with minimum cloud top temperatures $\sim -60^{\circ}\text{C}$ occurs above the Pyrénées between France and Spain and extends over $\sim 300 \times 300 \text{ km}^2$ in the early evening hours of 16 August 2013. The area above the thunderstorm is observed with a low-light video camera located in Clermont-Ferrand (45.8°N , 3.1°E , gray triangle). An intense positive lightning discharge with a peak current of $\sim 159 \text{ kA}$ occurs at 20:27:58.721 UTC (red cross) and causes a group of four sprites (dash-dotted lines) and a small luminosity column (white solid line). A consecutive negative lightning discharge with a peak current $\sim -6 \text{ kA}$ occurs at 20:27:58.755 UTC (blue dot). A radio receiver $\sim 77 \text{ km}$ northeast of the positive lightning discharge (gray square) records the electric field strength during the sequence of events.

[4] To determine relative luminosities against the bright moonlit background during the entire video sequence, we make use of morphological image processing. The first step of the image processing is to determine the gray level distribution function for each image pixel to identify the 10% largest luminosities for each pixel in the video sequence of 1.08 s length, which corresponds to 54 images. The identified candidate pixels are typically disconnected such that a spatial filter is applied to define a contiguous image area.

Here we make use of a morphological image opening with a 3×3 pixel large square as a structuring element. The image opening removes the isolated regions that are smaller than the structuring element while leaving the larger regions relatively unaffected [Gonzalez and Woods, 2008, pp. 635–639]. The resulting contiguous areas are used as an image mask to extract those pixels from the original image, which have unusual large luminosities. The relative luminosity of the sprite is determined here as the excess luminosity above the median, calculated with respect to time for each individual image pixel (Figure 2b).

[5] The image processing makes the body of the four sprites between $\sim 70 \text{ km}$ and $\sim 80 \text{ km}$ height, the diffuse glow between $\sim 80 \text{ km}$ and $\sim 90 \text{ km}$ height, and the tendrils extending down to $\sim 50 \text{ km}$ height immediately apparent. The sprite group is surrounded by several patches of luminosity with unknown significance. Yet one patch to the left of the sprite group (Figure 2b) rebrightens $\sim 40\text{--}60 \text{ ms}$ after the occurrence of the sprite group, and the rebrightening is followed by the development of a descending column of luminosity, i.e., a small column sprite, which is clearly significant (Figure 2c). The initial patch has a width of $\sim 1\text{--}2 \text{ km}$ and extends over $\sim 4\text{--}6 \text{ km}$ occupying $\sim 3\text{--}19 \text{ km}^3$ from ~ 80 to $\sim 85 \text{ km}$ height (Figure 2b), and it is located at least $\sim 15\text{--}20 \text{ km}$ to the left of the sprite group (Figure 1). This is a surprising result because the patch is completely disconnected from the diffuse glow above the sprite group. In other words, the electromagnetic field from the positive lightning discharge at 20:27:58.721 UTC was not sufficient to illuminate the area between the sprite group and the patch, but it was sufficient to illuminate the patch, which strongly indicates an irregularity of the ambient conditions in the volume of mesospheric air occupied by the patch. A weak negative lightning discharge with a peak current $\sim -6 \text{ kA}$ occurs at 20:27:58.755 UTC (Figure 1) and might have helped to initiate the rebrightening of the patch while other scattered lightning discharges occurring in the same area with similar peak currents (not shown in Figure 1) have no effect on the mesosphere above. The downward propagation speed during the rebrightening of the patch is estimated to be at least $\sim 1.25\text{--}2.50 \cdot 10^5 \text{ m s}^{-1}$ by using a vertical distance of $\sim 5 \text{ km}$ (5 pixels) traveled in $\sim 20\text{--}40 \text{ ms}$ (1–2 images). Note that this extremely slow downward descent of the column sprite

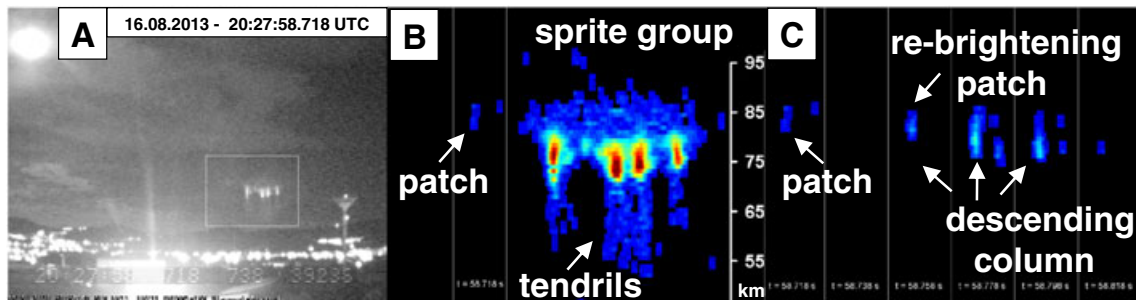


Figure 2. A descending column of luminosity appears after a group of four sprites recorded with the low-light video camera in Clermont-Ferrand. (a) The original video image shows a group of four sprites. Note that the sprite group was recorded against the full moon in the upper left corner of the video image. (b) The excess luminosity above the median of each pixel within the image mask exhibits a diffuse glow around $\sim 80\text{--}90 \text{ km}$ height above the body of the sprite group between ~ 70 and $\sim 80 \text{ km}$ and tendrils extending down to $\sim 50 \text{ km}$ height. A small patch of luminosity appears to the left of the sprite group between ~ 80 and $\sim 85 \text{ km}$ height. (c) This patch rebrightens at 20:27:58.758 and develops subsequently into a column which descends down to $\sim 74\text{--}76 \text{ km}$ height from 20:27:58.758 to 58.818.

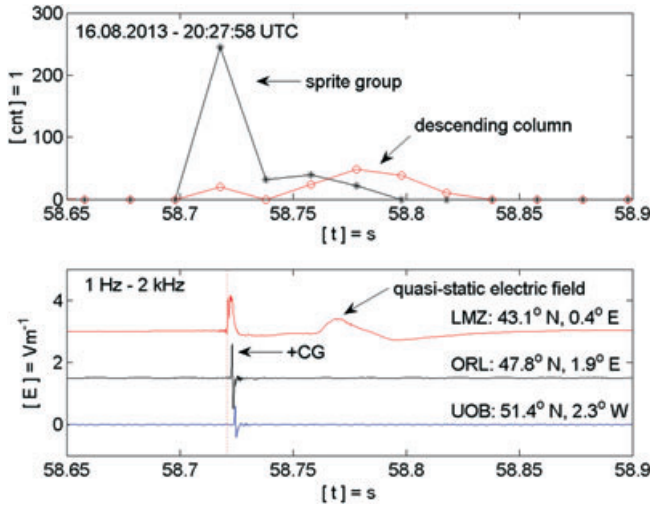


Figure 3. The descent of the column sprite is caused by a quasi-static electric field. (top) The relative luminosity of the descending column (red line) reaches its maximum ~ 60 ms after the peak luminosity of the sprite group (black line) as inferred from the video recordings. (bottom) The sprite group is caused by a positive cloud-to-ground discharge (+CG) at 20:27:58.7208 (red dotted line) as measured with three radio receivers at ~ 77 km (LMZ = Lannemezan, red line), ~ 604 km (ORL = Orléans, black line), and ~ 980 km (BTH = Bath, blue line) away from the lightning discharge in the frequency range from 1 Hz to 2 kHz. The nearest radio receiver ~ 77 km from the sprite shows an increase of the electric field strength, which causes the descending column of luminosity. The enhancement of the electric field strength for ~ 30 ms is not associated with electromagnetic radiation, which is detectable at the more distant radio receivers, and it is therefore attributed to a local quasi-static electric field. Note that the small time delays of the electromagnetic radiation from the +CG are caused by the wave propagation from the lightning discharge to the radio receivers.

is likely to require forcing by an additional electromagnetic field to occur.

[6] Vertical electric field strengths in the frequency range from ~ 1 Hz to ~ 400 kHz have been recorded with three wideband digital radio receivers near Lannemezan (43.1°N , 0.4°E), Orléans (47.8°N , 1.9°E), and Bath (51.4°N , 2.3°W) with a sampling frequency of 1 MHz, a resolution ~ 1 – 2 $\mu\text{V m}^{-1} \text{ Hz}^{-1/2}$, and a timing accuracy ~ 12 ns [Füllekrug, 2010]. The electric field recordings in the frequency range from 1 Hz to 2 kHz are compared to the relative luminosities of the sprite group and the patch (Figure 3). The relative luminosities are calculated from the square root of the integrated luminosities above the median for each image pixel within their respective image masks as inferred from the video recordings with 20 ms resolution. The nonlinear gain of the video camera and the saturation of the sprite group are not considered here because the comparison between the relative luminosity and the electric field recordings is entirely based on their temporal evolution. The relative luminosity of the sprite group coincides with the radio signal from the positive lightning discharge. The luminosity from the rebrightening patch and the descending column coincides with an electric field enhancement starting at 20:27:58.760 UTC (Figure 3), i.e., ~ 5 ms after the negative lightning

discharge with a peak current ~ -6 kA (Figure 1). The enhancement lasts for ~ 25 – 30 ms, and it is only seen by the radio receiver in Lannemezan ~ 77 km northeast of the positive lightning discharge. The radio receivers in Orléans and Bath are ~ 604 km and ~ 980 km away and do not exhibit the same electric field enhancement. It is therefore evident that this electric field enhancement results from a local quasi-static electric field which decays much faster with distance than the radiated electromagnetic field from the positive lightning discharge. The radiated field from the positive lightning discharge measured in Bath and Orléans indicates a charge moment change ~ 200 – 230 C km, which is inferred from the transfer function between the observed and calculated electric field spectrum [Füllekrug *et al.*, 2006, equations 1 and 2] using the short pulse approximation of the normal mode expansion with frequency-dependent ionospheric heights [Füllekrug and Rycroft, 2006, equation 2] from 1 Hz to 200 Hz. The observed charge moment change is relatively small because no lightning continuing current appears to be present in the electric field recordings.

[7] A high level of discharge activity is superimposed on the quasi-static electric field in Lannemezan (Figure 4, top left). The discharge activity is measured in the frequency range from 40 kHz to 400 kHz, and it exhibits numerous brief pulses which occur at a rate of ~ 10 – 20 per ms or ~ 1 – 2 per 100 μs . The intensity of the pulses and their occurrence rate are used to characterize the electric field variability ΔE of the discharge activity by calculating for each 100 μs long time interval the logarithmic mean

$$\mu = \frac{1}{N} \sum_k^N \log_{10} S(f_k) \quad (1)$$

where $S(f_k)$ are the spectral amplitude densities at $N = 37$ frequencies f_k from 40 kHz to 400 kHz separated by 10 kHz such that the electric field variability is finally given by $\Delta E = 10^\mu$. The agreement between the quasi-static electric field and the high-level discharge activity is excellent (Figure 4, top left and center left panels). A similar example of a high-level discharge activity superimposed on a quasi-static electric field (Figure 4, top right and center right panels) was observed in association with an extremely well-documented example of two consecutive lightning discharges [Füllekrug *et al.*, 2013b]. The source locations of the pulses occurring during the second lightning discharge have been mapped with an array of radio receivers [Mezentsev and Füllekrug, 2013] which shows that their origin are intracloud lightning discharges at an average height of ~ 3.5 km (Figure 4, bottom). The similarity of the two examples recorded near Lannemezan and Rustrel strongly suggests that the high level of discharge activity recorded in Lannemezan also results from intracloud lightning discharges.

3. Interpretation

[8] The interpretation of the experimental results is straightforward. The cloud-to-ground lightning discharge removes positive charge from the thundercloud, which causes the prompt appearance of the sprite group. A disconnected mesospheric irregularity concentrates the electric field lines from the positive cloud-to-ground lightning discharge. In this way, the mesospheric irregularity causes a localized enhancement of the electric field strength which

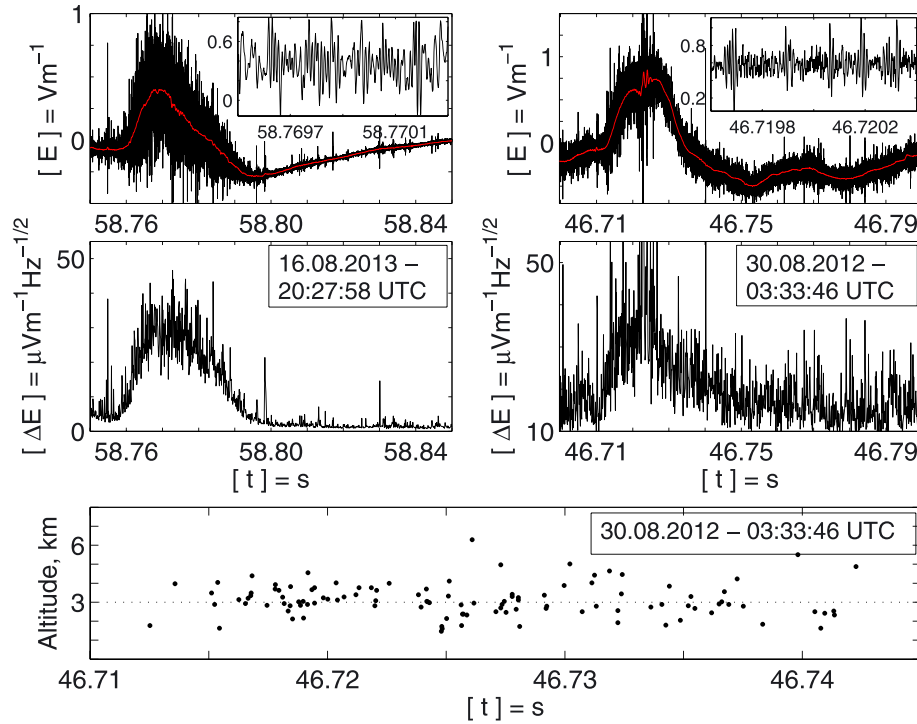


Figure 4. A high level of intracloud discharge activity is superimposed on the quasi-static electric field enhancement. (top) The high level of the discharge activity (black line) with $\sim 1\text{--}2$ pulses per $100\ \mu\text{s}$ (inset) is superimposed on the quasi-static electric field enhancement (red line). (top right) A similar example of a high-level discharge activity superimposed on a quasi-static electric field was recorded near Rustrel (43.9°N , 5.5°E) in southeastern France on 30 August 2012, 03:33:46 UTC. (middle) The quasi-static electric field is dissipated during the high-level discharge activity as measured by the electric field variability ΔE . (bottom) The high-level discharge activity occurs within the thundercloud at an altitude range $\sim 2\text{--}5$ km as mapped by a small-scale interferometric network of four radio receivers deployed near Rustrel.

results in the dielectric breakdown $\sim 15\text{--}20$ km away from the sprite group. This breakdown thereby illuminates the mesospheric irregularity, i.e., a luminous patch which occupies a volume $\sim 3\text{--}19\ \text{km}^3$. A weak negative lightning discharge initiates the rebrightening of the mesospheric irregularity, and the consecutive intracloud lighting discharges reorganize the charge distribution inside the thundercloud to compensate for the new charge configuration. This reorganization of charge results in a quasi-static electric field which causes the patch of mesospheric irregularity to develop into a columniform streamer which slowly descends at a speed of $\sim 1.25\text{--}2.50 \cdot 10^5\ \text{m s}^{-1}$. The charge moment Ql which causes the quasi-static electric field E can be inferred from the maximum-observed electric field $\sim 0.4\ \text{V m}^{-1}$ by use of the standard description for electrostatic fields

$$E = \frac{2Ql}{4\pi\epsilon_0(l^2 + d^2)^{3/2}} \quad (2)$$

[Uman, 1987, equation A.3, p. 315], where Q is the amount of charge at an average height l above the ground, ϵ_0 is the electric field constant, and d is the horizontal distance between the charge and the receiver. Assuming that the charge causing the quasi-static electric field is located near the mixed phase region of the thundercloud $\sim 4\text{--}6$ km height and that the horizontal distance is the same as the distance to the positive lightning discharge $\sim 77 \pm 20$ km results in a charge moment $\sim 4\text{--}20\ \text{C km}$. This is more than one order of magnitude below the charge moment $\sim 600\ \text{C km}$ which is typically required for the inception of sprites [Qin

et al., 2012a; Cummer and Lyons, 2005] and one order of magnitude below the charge moment $\sim 200\text{--}230\ \text{C km}$ observed here, which initiates the sprite group. However, this small charge moment $\sim 4\text{--}20\ \text{C km}$ does not reflect inception because it is only required to sustain the downward propagation of the columniform streamer which is also assisted by the radiated field from the high-level activity of intracloud discharges. Several other weak negative lightning discharges with small peak currents occur at scattered locations around the positive lightning discharge (not shown in Figure 1) without having any observable effect on the mesosphere above. As a result, the initial luminous patch is almost certainly produced by a localized concentration of electric field lines causing an enhancement of the electric field strength, which results in a small-scale dielectric breakdown of mesospheric air and thereby illuminates the mesospheric irregularity.

4. Discussion

[9] The excellent agreement between the quasi-static electric field enhancement and the intensity of the intracloud discharges does not offer any information on cause and effect. The chosen interpretation is based on the idea that the intracloud discharges drive the quasi-static electric field enhancement. However, the ongoing discussion on the initiation of lightning discharges during preliminary breakdown [Gurevich *et al.*, 2013; Gurevich and Karashtin, 2013; Dwyer, 2012, 2005; Gurevich *et al.*, 1992] is based on the

idea that streamer like microdischarges inside thunderclouds are caused by a preexisting quasi-static electric field. This line of thought would require a physical mechanism which causes the observed rise and decay of the quasi-static electric field. One possible mechanism could be for example electron pileup during which the positive cloud-to-ground discharge dislocates electrons with the heavier ions following later to neutralize the electric field produced by the dislocated electrons. Such kind of mechanism could be in operation, for example, if hydrometeors [Gurevich *et al.*, 2013; Gurevich and Karashtin, 2013] or aerosols [Füllekrug *et al.*, 2013, 2013b] contribute significantly to the observed quasi-static electric field.

[10] Quasi-static electric field measurements of sprites have previously been attributed to radiating current flowing in the body of the sprite [Hager *et al.*, 2012]. It is therefore interesting to estimate the charge per unit length, or line charge, of the descending column sprite because no long-range electromagnetic radiation is observed in this study. In this case, the maximum electric field is given by

$$E = \frac{2}{4\pi\epsilon_0} \int_{h_1}^{h_2} \frac{\rho z}{(z^2 + d^2)^{3/2}} dz = \frac{2\rho}{4\pi\epsilon_0} \left[-\frac{1}{\sqrt{z^2 + d^2}} \right]_{h_1}^{h_2} \quad (3)$$

[Uman, 1987, equation A.8, p. 317], where ρ is the average line charge along the length z extending from the bottom height $h_1 \approx 75$ km to the top height $h_2 \approx 85$ km of the columniform sprite at the time of maximum luminosity. The line charge stored on mesospheric aerosols or hydrometeors could then be as large as $\sim 30\text{--}60 \mu\text{C m}^{-1}$ assuming that the horizontal distance d to the columniform sprite is the same as the distance to the positive lightning discharge $\sim 77 \pm 20$ km. The inferred mesospheric line charge is one order of magnitude smaller than the line charge of a stratospheric leader stem in a gigantic jet [Neubert *et al.*, 2011], and it is hence indicative of a weak sprite. The observed quasi-static electric field could also result from a preceding halo [Bennett and Harrison, 2013; Parra-Rojas *et al.*, 2013; Kuo *et al.*, 2013], a bidirectional leader [Mallios *et al.*, 2013], or a resurgent lightning continuing current [Yashunin *et al.*, 2007]. However, all these processes would radiate electromagnetic waves to long-range distances which are not observed in this study.

[11] The theoretical modeling of sprite initiation suggests enhancements of the ambient electron density by one order of magnitude which possibly result from ionospheric disturbances by meteor trails, electrodynamic effects from thunderstorms and lightning, or gravity wave breaking [Qin *et al.*, 2013a; Kosar *et al.*, 2012] which are occasionally observed as irregular luminosities of elves and halos [van der Velde *et al.*, 2011] and might assist sprite initiation [Moudry *et al.*, 2003]. Electron density enhancements effectively concentrate the electric field lines and thereby increase the local electric field strength, which increases the likelihood for dielectric breakdown to occur. Electric field lines can also be concentrated at the surface of ambient air impurities such as solid particulates, e.g., iron from meteor ablation [Feng *et al.*, 2013; Gabrielli *et al.*, 2004; Rapp and Lübken, 2004] or ice crystals [Rapp and Thomas, 2006]. In either case, such kind of mesospheric irregularities near the mesopause would be illuminated by the electromagnetic field of lightning discharges as observed here.

[12] The detected mesospheric irregularity illuminates a maximum volume of $\sim 3\text{--}19 \text{ km}^3$ which is considered to be an upper bound as a result of the limited spatial and temporal resolution of the video images. For example, the initial luminous patch could result from a prompt column sprite with a typical width $\sim 0.2\text{--}0.6 \text{ km}$ [McHarg *et al.*, 2010] which propagates downward at a velocity $\sim 1\text{--}10 \cdot 10^6 \text{ m s}^{-1}$ [Li and Cummer, 2009] to reach a vertical extent of $\sim 4\text{--}6 \text{ km}$ within $\sim 0.5\text{--}5 \text{ ms}$ before it collapsed, such that the mesospheric irregularity would only account for some fraction of the observed luminosity. On the other hand, the horizontal extent of the patch $\sim 1\text{--}2 \text{ km}$ and the downward propagation speed $\sim 1.25\text{--}2.50 \cdot 10^5 \text{ m s}^{-1}$ during the rebrightening of the patch suggests that the mesospheric irregularity could occupy a volume up to $\sim 19 \text{ km}^3$ to account for all of the observed luminosity. In either case, the concentration of the electric field lines by the mesospheric irregularity at a significant distance from the sprite group effectively acts as an inception point for dielectric breakdown to occur, similar to the scattered inception points reported in association with a column and carrot sprite [Qin *et al.*, 2013b, Figure 1].

[13] Such mesospheric irregularities would be particularly effective in the absorption and scattering of radio waves with comparable wavelengths in the frequency range of hundreds of kilohertz which are used for the delivery of applications, e.g., atomic time transfer and marine navigation [Mezentsev and Füllekrug, 2013] remaining to be investigated in future studies.

5. Summary

[14] In summary, a mesospheric irregularity is illuminated by the electromagnetic pulse of an intense positive cloud-to-ground lightning discharge. A consecutive weak negative lightning discharge initiates the rebrightening of the irregularity which develops into a slowly descending columniform sprite assisted by a quasi-static electric field enhancement which is driven by a high level of intracloud lightning discharge activity. Future studies on the illumination of mesospheric irregularities by lightning discharges could make use of video cameras with a smaller noise level of the recording chip, sampling at higher frequencies, and use telescopic imaging to focus around the mesopause region. Such kind of experimental scientific program might assist an improved understanding of the absorption and scattering of 100 kHz radio waves which are used for the delivery of applications such as atomic time transfer and marine navigation.

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